Search for $\Omega(2012) \rightarrow K\Xi(1530) \rightarrow K\pi\Xi$ at Belle


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I. INTRODUCTION

Very recently a new state, the excited $\Omega(2012)$ baryon, has been observed by the Belle collaboration [1] in the $\Xi\Lambda$ invariant mass spectra using data samples collected at the $\Upsilon(1S, 2S, 3S)$ energies, with measured mass $M = [2012.4 \pm 0.7$ (stat) $\pm 0.6$ (syst)] MeV/c$^2$ and width $\Gamma = [6.4 \pm 2.5$ (stat) $\pm 1.6$ (syst)] MeV. The observed spacing in the $\Omega$ mass spectrum between the ground state and this excited state ($\sim$340 MeV/c$^2$) is smaller than that for other $\Omega^-$ excited states [2], and is more similar to the negative-parity orbital excitations of many other baryon pairs such as $\Lambda$ and $\Lambda(1405)$ or $\Lambda^+_c$ and $\Lambda^+_b$(2595).

After the initial observation of the $\Omega(2012)$, several theoretical interpretations of that state were offered [3–11]. Although it is generally accepted that $\Omega(2012)$ is a $1^P$ orbital excitation of the ground-state $\Omega$ baryon with quark content $sss$ and quantum numbers $J^P = \frac{1}{2}^-$, Refs. [7–10] propose an alternative interpretation as a $K\Xi(1530)$ hadronic molecule. These models predict a large decay width for $\Omega(2012) \to K\pi\Xi$. In Ref. [7], the decay $\Omega(2012) \to K\pi\Xi$ is predicted to dominate over $\Omega(2012) \to K\Xi$, while in Refs. [8–10], the production rates of the $\Omega(2012)$ are almost similar in $K\pi\Xi$ and $K\Xi$ decay channels. The authors in Ref. [11] also discuss the three-body decay of the $\Omega(2012)$. They regard $\Omega(2012)$ as a member of the compact decuplet states only if the sum of branching fractions of the $\Omega(2012) \to K\pi\Xi$ and $\Omega(2012) \to \Omega\pi\Xi$ is not too large ($< 70\%$).

In this paper, we report on a search for $\Omega(2012) \to K\Xi(1530) \to K\pi\Xi$ using $\Upsilon(1S, 2S, 3S)$ data samples collected by the Belle experiment at the KEKB asymmetric-energy $e^+e^-$ collider [12,13]. Note that charge-conjugate modes are implied throughout, unless explicitly stated otherwise.

II. THE DATA SAMPLE AND BELLE DETECTOR

The Belle data used in this analysis correspond to 5.7 fb$^{-1}$ of integrated luminosity at the $\Upsilon(1S)$ resonance, 24.9 fb$^{-1}$ at the $\Upsilon(2S)$ resonance, and 2.9 fb$^{-1}$ at the $\Upsilon(3S)$ resonance. The Belle detector [14,15] is a large solid-angle magnetic spectrometer consisting of a silicon vertex detector (SVD), a 50-layer central drift chamber (CDC), an array of aerogel threshold Cherenkov counters (ACC), a barrel arrangement of time-of-flight scintillation counters (TOF), and an electromagnetic calorimeter comprised of CsI(Tl) crystals (ECL) located inside a superconducting solenoid coil providing a 1.5 T magnetic field. An iron flux-return yoke instrumented with resistive plate chambers (KLM) located outside the coil is used to detect $K_L^0$ mesons and to identify muons.

Large signal Monte Carlo (MC) samples (1 million events for each studied process) are generated using the EVTGEN [16] code to simulate the expected signal event topology and estimate the signal detection efficiency. The processes $\Upsilon(1S, 2S, 3S) \to \Omega(2012) \to K\Xi(1530) + \text{anything} \to K\Xi + \text{anything}$ are simulated; the mass and width of $\Omega(2012)$ are fixed at 2012 GeV/c$^2$ and 6.4 MeV [1], respectively. To assess possible backgrounds arising from the continuum ($e^+e^- \to q\bar{q}$ with $q = u, d, s, c$), we generate such events at center-of-mass energies of $\Upsilon(1S), \Upsilon(2S), \text{and } \Upsilon(3S)$ resonances using the Lund fragmentation model in PYTHIA [17]. Inclusive $\Upsilon(1S)$ and $\Upsilon(2S)$ MC samples, corresponding to four times the
ACC is used to identify charged kaons and pions based on their momenta: more than 0.8, where $L_K$ and $L_\pi$ are the likelihood values for the kaon and pion hypotheses, respectively. Tracks with $R_K = L_K/(L_K + L_\pi) < 0.4$ are identified as pions with an efficiency of 96%, while 8% of kaons are misidentified as pions; tracks with $R_K > 0.6$ are identified as kaons with an efficiency of 95%, while 6% of pions are misidentified as kaons.

An ECL cluster is treated as a photon candidate if it does not match the extrapolation of any charged track reconstructed by the tracking systems (CDC and SVD) into the calorimeter. The $\pi^0$ candidates are reconstructed from two photons having energy exceeding 50 MeV in the barrel or 100 MeV in the endcaps. To avoid contamination from neutral hadrons, we reject neutral showers if the ratio of the energy deposited in the central array of ECL cells to that deposited in the surrounding array of $5 \times 5$ cells is less than 0.8. The $\pi^0 \rightarrow \gamma\gamma$ candidates are also required to have an energy balance parameter $|E_1 - E_2|/(E_1 + E_2)$ smaller than 0.8, where $E_1$ ($E_2$) is the energy of the first (second) photon in the laboratory frame. To further reduce the combinatorial background, the momentum of the $\pi^0$ candidate is required to exceed 200 MeV/c. We define the $\pi^0$ signal region as $|M_{\gamma\gamma} - m_{\pi^0}| < 12$ MeV/c$^2$ (~2$\sigma$), where $m_{\pi^0}$ is the $\pi^0$ nominal mass [2]. For each selected $\pi^0$ candidate a mass-constrained fit is performed to improve its momentum resolution.

The $K_S^0$ candidates are reconstructed via the $K_S^0 \rightarrow \pi^+\pi^-$ decay, and the identification is enhanced by selecting on the outputs of a neural network [18]. The network uses the following input variables [19]: the $K_S^0$ momentum in the lab frame, the distance along the $z$ axis between the two track helices at their closest approach, the $K_S^0$ flight length in the $r - \phi$ plane, the angle between the $K_S^0$ momentum and the vector joining the interaction point (IP) to the $K_S^0$ decay vertex, the angle between the pion momentum and the lab frame direction in the $K_S^0$ rest frame, the distances of closest approach in the $r - \phi$ plane between the IP and the two pion helices, the number of hits in the CDC for each pion track, and the presence or absence of hits in the SVD for each pion track.

Candidate $\Lambda$ decays are reconstructed from $p\pi^-$ pairs with a production vertex significantly separated from the IP. For the $\Xi^-(\rightarrow \Lambda\pi^-)$ and $\Xi^0(\rightarrow \Lambda\pi^0)$ candidates, the vertex fits are performed and the positive $\Xi^-$ and $\Xi^0$ flight distances are required. The selected $\Xi^-(\rightarrow \Lambda\pi^-)$ and $\Xi^0(\rightarrow \Lambda\pi^0)$ candidates are the same as those in Ref. [1]. The $\Xi^-$ candidates are the likelihood values for the kaon and pion hypotheses, respectively. Tracks with $R_K = L_K/(L_K + L_\pi) < 0.4$ are identified as pions with an efficiency of 96%, while 8% of kaons are misidentified as pions; tracks with $R_K > 0.6$ are identified as kaons with an efficiency of 95%, while 6% of pions are misidentified as kaons.

![Figure 1](https://example.com/figure1.png)

**FIG. 1.** Distributions of (a) $M(\Xi(1530)^0K^-)$ versus $M(\Xi^-\pi^+)$, (a) $M(\Xi(1530)^-K^0_S)$ versus $M(\Xi^-\pi^0)$, (a) $M(\Xi(1530)^-K^0_S)$ versus $M(\Xi^-\pi^0)$, and (d) $M(\Xi(1530)^0K^-)$ versus $M(\Xi^0\pi^0)$ from signal MC samples. The dotted lines bound the $\Xi(1530)$ signal region.
and $\Xi^0$ are kinematically constrained to their nominal masses [2], and then combined with a $\pi^\pm$ or $\pi^0$ to form a $\Xi(1530)^-$ or $\Xi(1530)^0$ candidate. Finally, the selected $\Xi(1530)$ candidate is combined with a $K^-$ or $K^0_S$ to form the $\Omega(2012)$ candidate. In this last step, a vertex fit is performed for the $K\pi\Xi$ final state to improve the momentum resolutions and suppress the backgrounds, requiring $\chi^2_{\text{vertex}} < 20$, corresponding to an estimated selection efficiency exceeding 95%. Reconstruction spans the $\Omega(2012)^-\rightarrow \Xi^-\pi^-K^-$, $\Xi^-\pi^0K^0_S$, $\Xi^-\pi^-K^0_S$, and $\Xi^0\pi^0K^-$ three-body decay modes of $\Omega(2012)$.

Before searching for $\Omega(2012)\rightarrow K\Xi(1530)\rightarrow K\pi\Xi$, a cross-check on the previously reconstructed $\Omega(2012)\rightarrow \Xi\pi\Xi$ candidates uses well-reconstructed tracks, particle identifications, and vertex fitting technique in a way similar to the methods in Ref. [1]. As a result, the signal yields from the simultaneous fit of the $\Omega(2012)^-\rightarrow \Xi^-\pi^-K^-$ and $\chi^2_{\text{vertex}} < 20$. The obtained mass and width for the $\Omega(2012)$ are $M = (2012.1 \pm 0.7)$ MeV/$c^2$ and $\Gamma = (6.9^{+2.5}_{-2.0})$ MeV, where the uncertainties are statistical only. Our results are consistent with those in Ref. [1] within errors.

**B. The distributions from signal MC samples**

After all event selection requirements, Fig. 1 shows the distributions of the $\Xi\pi K$ invariant mass versus the $\Xi\pi$ invariant mass from signal MC samples. Due to phase space limitations, events at high $\Xi\pi$ and/or low $\Xi\pi K$ mass are kinematically forbidden. We define the optimized $\Xi(1530)$ signal region as $1.49 \text{ GeV}/c^2 < M(\Xi\pi) < 1.53 \text{ GeV}/c^2$ (discussed below), between the blue dashed lines in Fig. 1.

The invariant mass distributions from MC signal simulations of $\Xi(1530)^0\rightarrow \Xi^-\pi^-/\Xi^0\pi^0)K^-$ and $\Xi(1530)^-\rightarrow \Xi^-\pi^0/\Xi^0\pi^-)K^0_S$ are shown in Fig. 2. The signal shape of the $\Omega(2012)$ is described by a Breit-Wigner (BW) function convolved with a Gaussian function, where the BW mass and width are fixed to 2.0124 GeV/$c^2$ and 6.4 MeV [1], respectively, and the mass-resolution, i.e., Gaussian width is determined in the fit.

**C. $\Xi(1530)$ signals in $\Upsilon(1S.2S.3S)$ data**

After imposing our selection criteria, the invariant mass spectra of $\Xi(1530)^0\rightarrow \Xi^-\pi^+.\Xi^0\pi^0$, and $\Xi(1530)^-\rightarrow \Xi^-\pi^0$, $\Xi^0\pi^- $ candidates are shown in Figs. 3(a)–3(d). Clear signals of $\Xi(1530)^0$ and $\Xi(1530)^-$ are observed in the modes $\Xi(1530)^0\rightarrow \Xi^-\pi^+$ and $\Xi(1530)^-\rightarrow \Xi^-\pi^0$, $\Xi^0\pi^- $.

We fit all the invariant mass distributions, modeling the $\Xi(1530)$ peaks with the convolution of a BW and a Gaussian function and the background as a second-order polynomial. In the fits, the BW parameters are unconstrained, while the Gaussian widths are fixed according to MC simulations. The fit values are consistent with the world averages within their respective errors [2]. For $\Xi(1530)^0\rightarrow \Xi^0\pi^0$, the mass and width of $\Xi(1530)^0$ are
fixed to the Particle Data Group (PDG) values [2] since the signal is not clear due to large combinatorial backgrounds. The results of the fits are listed in Table I.

**D. \(\Xi(2012) \rightarrow \Xi \pi K\) mass distributions in \(\Upsilon(1S, 2S, 3S)\) data**

Considering phase space limitations and our finite mass resolution, we require 1.49 GeV/c^2 < \(M(\Xi \pi)\) < 1.53 GeV/c^2 to select \(\Xi(1530)\) signals as efficiently as possible, as indicated by the red arrows in Fig. 3. We optimize this requirement by maximizing the figure of merit \(N_{\text{sig}}/\sqrt{N_{\text{sig}} + N_{\text{bkg}}}\) value with the mode \(\Omega(2012)^- \rightarrow \Xi(1530)^0(\rightarrow \Xi^+ \pi^-)K^-\), where \(N_{\text{sig}}\) is number of fitted signal events in the signal MC sample assuming \(B(\Upsilon(1S, 2S, 3S) \rightarrow \Omega(2012)^- + \text{anything}) \times B(\Omega(2012)^- \rightarrow \Xi(1530)^0K^-) = 10^{-6}\) and \(N_{\text{bkg}}\) is the number of estimated background events in the \(\Omega(2012)^-\) signal region using inclusive MC samples. The candidate signal region for the \(\Xi(1530)\) coincides with the predicted mass interval from Ref. [9].

After application of the above selection criteria, Fig. 4 shows the invariant mass distributions of \(\Xi(1530)^0(\rightarrow \Xi^- \pi^+/\Xi^0 \pi^0)K^-\) and \(\Xi(1530)^-\rightarrow(\rightarrow \Xi^0 \pi^-/\Xi^0 \pi^0)K_S^0\). From these distributions, no obvious \(\Omega(2012)^-\) signal is observed. The shapes of the \(\Omega(2012)\) signals in the fits are described by BW functions convolved with Gaussian resolution functions; the background shapes are described by a threshold function. The parameters of the BW functions are fixed to the mass and width of the \(\Omega(2012)\) [1], and the mass resolutions are fixed to those from fits to signal MC samples (1.5, 2.6, 1.7, and 2.8 MeV for the \(\Omega(2012)^-\rightarrow\Xi(1530)^0(\rightarrow \Xi^- \pi^+/\Xi^0 \pi^0)K^-\), \(\Xi(1530)^-\rightarrow(\rightarrow \Xi^0 \pi^-)K_S^0\), \(\Xi(1530)^-\rightarrow(\rightarrow \Xi^0 \pi^0)K_S^0\), and \(\Xi(1530)^0(\rightarrow \Xi^0 \pi^0)K^-\) decay modes, respectively). The threshold function has the form \((M(\Xi K \pi) - \chi)\), \(\exp(c_1(M(\Xi K \pi) - \chi) + c_2(M(\Xi K \pi) - \chi)^2),\)

**TABLE I.** Mass resolution from MC simulations, and the mass and width for inclusive \(\Xi(1530)\) signals from fits to the Belle data.

<table>
<thead>
<tr>
<th>Mode</th>
<th>Resolution ((\text{MeV/}c^2))</th>
<th>Mass ((\text{MeV/}c^2))</th>
<th>Width ((\text{MeV}))</th>
</tr>
</thead>
<tbody>
<tr>
<td>(\Xi(1530)^0 \rightarrow \Xi^- \pi^+)</td>
<td>2.34 ± 0.14</td>
<td>1532.47 ± 0.03</td>
<td>9.0 ± 0.3</td>
</tr>
<tr>
<td>(\Xi(1530)^- \rightarrow \Xi^- \pi^0)</td>
<td>2.96 ± 0.17</td>
<td>1535.07 ± 0.37</td>
<td>12.9 ± 1.8</td>
</tr>
<tr>
<td>(\Xi(1530)^- \rightarrow \Xi^0 \pi^-)</td>
<td>2.44 ± 0.15</td>
<td>1535.11 ± 0.09</td>
<td>10.6 ± 0.2</td>
</tr>
<tr>
<td>(\Xi(1530)^0 \rightarrow \Xi^0 \pi^0)</td>
<td>4.14 ± 0.26</td>
<td>1531.80 (PDG value)</td>
<td>9.1 (PDG value)</td>
</tr>
</tbody>
</table>
where the parameters $\alpha$, $c_1$, and $c_2$ are free; the threshold parameter $x$ is fixed at 1.97 GeV/c$^2$ from the MC simulations. The yields of $\Omega(2012)$ signal events from the unbinned extended maximum-likelihood fits are obtained; they are listed in Table II, together with the reconstruction efficiency, signal significance, and the upper limit at 90% credibility level [20] (C.L.) on the signal yield for each $\Omega(2012)$ decay mode. In addition, no peaking backgrounds are found from the inclusive MC samples.

### E. The ratios of the branching fractions for $\Omega(2012) \rightarrow K\pi\Xi$ relative to $K\Xi$

We define the ratios $\mathcal{R}_{\Xi^0}^{\Xi^-} = \frac{\mathcal{B}(\Omega(2012) \rightarrow \Xi(1530)^0(\rightarrow \Xi^-\pi^+)K^-)}{\mathcal{B}(\Omega(2012) \rightarrow \Xi^-K^0)}$, $\mathcal{R}_{\Xi^0}^{\Xi^0} = \frac{\mathcal{B}(\Omega(2012) \rightarrow \Xi(1530)^0(\rightarrow \Xi^0\pi^+)K^-)}{\mathcal{B}(\Omega(2012) \rightarrow \Xi^-K^0)}$, $\mathcal{R}_{\Xi^0}^{\Xi^-} = \frac{\mathcal{B}(\Omega(2012) \rightarrow \Xi(1530)^0(\rightarrow \Xi^-\pi^+)K^-)}{\mathcal{B}(\Omega(2012) \rightarrow \Xi^-K^0)}$, $\mathcal{R}_{\Xi^0}^{\Xi^-} = \frac{\mathcal{B}(\Omega(2012) \rightarrow \Xi(1530)^0(\rightarrow \Xi^-\pi^+)K^-)}{\mathcal{B}(\Omega(2012) \rightarrow \Xi^-K^0)}$, $\mathcal{R}_{\Xi^0}^{\Xi^0} = \frac{\mathcal{B}(\Omega(2012) \rightarrow \Xi(1530)^0(\rightarrow \Xi^-\pi^+)K^-)}{\mathcal{B}(\Omega(2012) \rightarrow \Xi^-K^0)}$, and $\mathcal{R}_{\Xi^0}^{\Xi^-} = \frac{\mathcal{B}(\Omega(2012) \rightarrow \Xi(1530)^0(\rightarrow \Xi^-\pi^+)K^-)}{\mathcal{B}(\Omega(2012) \rightarrow \Xi^-K^0)}$ and determine their values as follows:

### TABLE II. The reconstruction efficiency ($\epsilon$), signal significance ($\sigma$), signal yield ($N_{fit}$), and the upper limit at 90% C.L. ($N_{UL}$) on the signal yield for each $\Omega(2012)$ decay mode.

<table>
<thead>
<tr>
<th>Mode</th>
<th>$\epsilon$ (%)</th>
<th>$\sigma$</th>
<th>$N_{fit}$</th>
<th>$N_{UL}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$\Omega(2012)^+ \rightarrow \Xi(1530)^0(\rightarrow \Xi^-\pi^+)K^-$</td>
<td>8.71 ± 0.06</td>
<td>1.8</td>
<td>22.5 ± 12.9</td>
<td>41.0</td>
</tr>
<tr>
<td>$\Omega(2012)^+ \rightarrow \Xi(1530)^0(\rightarrow \Xi^-\pi^+)K^0$</td>
<td>1.26 ± 0.01</td>
<td>...</td>
<td>-3.5 ± 11.6</td>
<td>16.6</td>
</tr>
<tr>
<td>$\Omega(2012)^+ \rightarrow \Xi(1530)^0(\rightarrow \Xi^-\pi^+)K^-_S$</td>
<td>2.06 ± 0.02</td>
<td>...</td>
<td>-1.0 ± 3.6</td>
<td>7.2</td>
</tr>
<tr>
<td>$\Omega(2012)^+ \rightarrow \Xi(1530)^0(\rightarrow \Xi^-\pi^+)K^-$</td>
<td>0.75 ± 0.01</td>
<td>...</td>
<td>-12.0 ± 9.8</td>
<td>13.2</td>
</tr>
</tbody>
</table>
Table III. The values of the $R^\Xi_{K^0}$, $R^\Xi_0K^0$, $R^\Xi_0K^-$, $R^\Xi_0K^-\pi^-$, $R^\Xi_0K^-\pi^0$, and $R^\Xi_0K^-\pi^+$.

<table>
<thead>
<tr>
<th>The ratio</th>
<th>The value</th>
</tr>
</thead>
<tbody>
<tr>
<td>$R^\Xi_{K^0}$</td>
<td>(5.0 ± 2.9)%</td>
</tr>
<tr>
<td>$R^\Xi_0K^0$</td>
<td>(−15.8 ± 52.3)%</td>
</tr>
<tr>
<td>$R^\Xi_0K^-$</td>
<td>(−2.3 ± 8.4)%</td>
</tr>
<tr>
<td>$R^\Xi_0K^-\pi^-$</td>
<td>(−26.8 ± 21.9)%</td>
</tr>
<tr>
<td>$R^\Xi_0K^-\pi^0$</td>
<td>(4.2 ± 2.5)%</td>
</tr>
<tr>
<td>$R^\Xi_0K^-\pi^+$</td>
<td>(−2.8 ± 10.0)%</td>
</tr>
</tbody>
</table>

The ratio

$$R^\Xi_{K^0} = \frac{B(\Omega(2012) \to \Xi(1530)\rightarrow \Xi^0n^0)K^-)}{B(\Omega(2012) \to \Xi^-)}$$

$$= \frac{N_{fit}^\Xi \times e_f}{N_0^\Xi \times e_0 \times B(\Xi^0 \to \gamma\gamma)}$$

$$R^\Xi_0K^0 = \frac{B(\Omega(2012) \to \Xi(1530)\rightarrow \Xi^-\pi^+)K^-)}{B(\Omega(2012) \to \Xi^-)}$$

$$= \frac{N_{fit}^\Xi \times e_f \times B(\Xi^0 \to \Lambda\pi^+)}{N_0^\Xi \times e_0 \times B(\Xi^- \to \Xi^-)}$$

$$R^\Xi_0K^-\pi^- = \frac{B(\Omega(2012) \to \Xi(1530)\rightarrow \Xi^-\pi^-)K^-)}{B(\Omega(2012) \to \Xi^-)}$$

$$= \frac{N_{fit}^\Xi \times e_f \times B(\Xi^- \to \Lambda\pi^-)}{N_0^\Xi \times e_0 \times B(\Xi^- \to \Xi^-)}$$

where the errors are statistical only; $N_{fit}^\Xi$, $N_0^\Xi$, $N_{fit}^\Xi$, $N_0^\Xi$, $N_{fit}^\Xi$, $N_0^\Xi$, and $N_{fit}^\Xi$ are the fitted signal yields in the modes $\Omega(2012)\rightarrow \Xi(1530)\rightarrow \Xi^-\pi^+K^-$, $\Xi(1530)\rightarrow \Xi^-\pi^-K^0$, $\Xi(1530)\rightarrow \Xi^-\pi^0K^0$, $\Xi(1530)\rightarrow \Xi^-\pi^+K^0$, $\Xi(1530)\rightarrow \Xi^-\pi^-K^0$, $\Xi(1530)\rightarrow \Xi^-\pi^0K^0$, respectively; $e_1$, $e_2$, $e_3$, $e_4$, $e_5$, and $e_6$ are the corresponding efficiencies for each mode. The values of $N_{fit}^\Xi$, $N_0^\Xi$, $N_{fit}^\Xi$, $N_0^\Xi$, $N_{fit}^\Xi$, $N_0^\Xi$, and $N_{fit}^\Xi$ are listed in Table II. The values of $N_0^\Xi$, $N_0^\Xi$, and $e_6$ are 279 ± 71, 242 ± 48, (15.7 ± 0.2)% and (4.0 ± 0.1)%. In our calculations, we use the standard value of $B(\Xi^0 \to K^0) = 0.35$.

Finally, the values of $R^\Xi_{K^0}$, $R^\Xi_0K^0$, $R^\Xi_0K^-$, $R^\Xi_0K^-\pi^-$, $R^\Xi_0K^-\pi^0$, and $R^\Xi_0K^-\pi^+$ are obtained; they are listed in Table III.

F. Simultaneous fit results

Considering the branching fractions of $\Omega(2012)\rightarrow \Xi^-K^0$ and $\Omega(2012)\rightarrow \Xi^0K^-$ and the ratios of branching fractions of the three-body decay modes of $\Omega(2012)$ are known, the ratio of expected signal yields between each $\Omega(2012)$ three-body decay mode can be calculated. With such constraints, we perform a simultaneous fit to obtain the upper limit on $R^\Xi_{E_K} = B(\Omega(2012) \to \Xi(1530)\rightarrow \Xi^-\pi^+)K^-/B(\Omega(2012) \to \Xi^-)$.

Taking $B(\Omega(2012)\rightarrow \Xi^-\pi^+)K^-/B(\Omega(2012)\rightarrow \Xi^0K^-) = 1.0:1.2$ [1] and $B(\Omega(2012)\rightarrow \Xi(1530)\rightarrow \Xi^-\pi^+K^-)/B(\Omega(2012)\rightarrow \Xi(1530)\rightarrow \Xi^-\pi^0K^-) = 1:1$ according to isospin symmetry, we derive that $R^\Xi_{E_K}:R^\Xi_{E_K}:R^\Xi_{E_K}:R^\Xi_{E_K} = 1:1:1:1$. Thus, according to Eqs. (1)–(4), we have:

$$N_{fit}^{\Xi}/N_0^{\Xi}/N_3^{\Xi}/N_{fit}^{\Xi} = 87.2%:2.2%:7.0%:3.6%.$$ (7)

An unbinned extended maximum-likelihood simultaneous fit to all three-body decay modes is now performed. In the simultaneous fit, the ratios of the expected observed $\Omega(2012)$ signals between each decay channel are fixed according to Eq. (7). The functions used to describe the signal and background shapes are parametrized as before. The fit result is shown in Fig. 5 from the combined $\Upsilon(1S,2S,3S)$ data samples, corresponding to a total fit yield of 22.4 ± 14.0. The statistical significance of the $\Omega(2012)$ signal is 1.6σ. Finally, we determine

$$R^\Xi_{E_K} = \frac{B(\Omega(2012) \to \Xi(1530)\rightarrow \Xi^-\pi^+)K^-)}{B(\Omega(2012) \to \Xi^-)}$$

$$= (6.0 \pm 3.7(stat) \pm 1.3(sys))%.$$ (8)
G. Systematic uncertainties

We now discuss the systematic uncertainties inherent in our measurements of the ratios $R^{\Xi^-\pi^+\bar{K}^-}/R^{\Xi^-\pi^+\bar{K}^0}$, $R^{\Xi^-\pi^-\bar{K}^0}/R^{\Xi^-\pi^-\bar{K}^0}$, $R^{\Xi^-\pi^+\bar{K}^-}/R^{\Xi^-\pi^+\bar{K}^0}$, $R^{\Xi^-\pi^-\bar{K}^-}/R^{\Xi^-\pi^-\bar{K}^0}$, $R^{\Xi^-\pi^-\bar{K}^-}/R^{\Xi^-\pi^-\bar{K}^0}$, $R^{\Xi^-\pi^-\bar{K}^-}/R^{\Xi^-\pi^-\bar{K}^0}$, and $R^{\Xi^-\pi^+\bar{K}^-}/R^{\Xi^-\pi^+\bar{K}^0}$. These include detection efficiency (tracking efficiency, kaon and pion particle ID, $\Lambda$, $K^0_s$, and $\pi^0$ reconstruction), the statistical error in the MC efficiency, the branching fractions of possible intermediate states, the $\Omega(2012)$ resonance parameters, any possible bias in reconstructed mass (as evaluated from the difference between the reconstructed $\Xi^0$ mass and the world average value), as well as the overall fit uncertainty.

Based on a study of $D^{*+} \to D^0(\to K^0_s\pi^+\pi^-)\pi^+$, the uncertainty in tracking efficiency is taken to be 0.35% per track. The uncertainties in particle identification are studied via a low-background sample of $D^+$ decay for charged kaons and pions. The studies show uncertainties of 1.3% for each charged kaon and 1.1% for each charged pion. The uncertainty in $\Lambda$ selection is 3% [21]. Differences in $K^0_s$ selection efficiency determined from data and MC simulations give a relation of $1 - \epsilon_{\text{data}}/\epsilon_{\text{MC}} = (1.4 \pm 0.3)\%$ [22]; 1.7% is taken as a conservative systematic uncertainty. For $\pi^0$ reconstruction, the efficiency correction and systematic uncertainty are estimated from a sample of $\tau^- \to \pi^-\pi^0\nu$. We find a 2.25% systematic uncertainty on $\pi^0$ reconstruction efficiency. In the measurements of $R^{\Xi^-\pi^+\bar{K}^-}/R^{\Xi^-\pi^+\bar{K}^0}$, $R^{\Xi^-\pi^-\bar{K}^-}/R^{\Xi^-\pi^-\bar{K}^0}$, $R^{\Xi^-\pi^+\bar{K}^-}/R^{\Xi^-\pi^+\bar{K}^0}$, and $R^{\Xi^-\pi^-\bar{K}^-}/R^{\Xi^-\pi^-\bar{K}^0}$, the common sources of systematic uncertainties such as $\Xi$ selection cancel; the individual errors are summed in quadrature to obtain the total detection efficiency uncertainty. For the measurement of $R^{\Xi^-\pi^+\bar{K}^-}/R^{\Xi^-\pi^+\bar{K}^0}$, to determine the total detection efficiency, the systematic errors for each final state and the errors from tracking, particle identification, $\Lambda$, $K^0_s$, and $\pi^0$ reconstruction are first summed in quadrature to obtain $\sigma_i$. Then, the total systematic uncertainty for detection efficiency ($\sigma_{\text{DE}}$) is determined using standard error propagation as follows:

$$\sigma_{\text{DE}} = \sqrt{\frac{\Sigma_i(W_i \times \sigma_i)^2 + \Sigma_j(W_j \times \sigma_j)^2}{(\Sigma_i W_i)^2}} = 7.3\%$$

Here, $W_i$ is the weight factor of the branching fraction in the $i$th ($i = 0, 1, 2, 3$) mode of the $\Omega(2012) \to \Xi\pi\bar{K}$ decay; $W_j$ ($j = 0, 1$) is the relative weight for the $j$th mode of $\Omega(2012) \to \Xi\pi\bar{K}$ decay.

The statistical uncertainty in the determination of the efficiency from MC simulations is less than 1.0%. In the calculation of $R^{\Xi^-\pi^+\bar{K}^-}/R^{\Xi^-\pi^+\bar{K}^0}$, only the branching fractions of intermediate states $B(K^0_s \to \pi^+\pi^-)$ and $B(\pi^0 \to \gamma\gamma)$ are included; the corresponding uncertainties are 0.072% and 0.035% [2], respectively, which are sufficiently small to be neglected. The uncertainty in the $\Omega(2012)$ resonance parameters is estimated by toggling the values of resonance mass and width by $\pm 1\sigma$ and refitting. The largest differences compared to the nominal fit results are taken as the systematic uncertainties associated with the $\Omega(2012)$ resonance parameters. The uncertainty in the $\Xi^0$ mass is estimated by comparing the numbers of the signal yields of the $\Omega(2012)$ for the case where the mass of the reconstructed $\Xi^0$ is fixed at the found peak value versus the case where the mass is fixed to the nominal mass [2]. According to the $\Xi(1530)\bar{K}$ invariant mass distributions in inclusive MC samples, we find that the threshold mass value falls within the $[1.96, 1.98]$ GeV/$c^2$ interval. The systematic error in the background parametrization is estimated by comparing the yields when the threshold mass is changed by $\pm 10$ MeV/$c^2$ relative to the nominal fit (for which the threshold is fixed at $1.97$ GeV/$c^2$). By extending the fitted region of the $M(\Xi\pi\bar{K})$, e.g., 2.2 to 2.3 GeV/$c^2$, the upper limits at 90% C.L. on the $\Omega(2012)$ signal yields are not changed. Such systematic uncertainty due to the fit region can be neglected.

All the uncertainties are summarized in Table IV, and, assuming all errors are independent, summed in quadrature to give the total systematic uncertainty.

<table>
<thead>
<tr>
<th>Source</th>
<th>$R^{\Xi^-\pi^+\bar{K}^-}/R^{\Xi^-\pi^+\bar{K}^0}$</th>
<th>$R^{\Xi^-\pi^-\bar{K}^-}/R^{\Xi^-\pi^-\bar{K}^0}$</th>
<th>$R^{\Xi^-\pi^+\bar{K}^-}/R^{\Xi^-\pi^+\bar{K}^0}$</th>
<th>$R^{\Xi^-\pi^-\bar{K}^-}/R^{\Xi^-\pi^-\bar{K}^0}$</th>
<th>$R^{\Xi^-\pi^+\bar{K}^-}/R^{\Xi^-\pi^+\bar{K}^0}$</th>
<th>$R^{\Xi^-\pi^-\bar{K}^-}/R^{\Xi^-\pi^-\bar{K}^0}$</th>
<th>$R^{\Xi^-\pi^+\bar{K}^-}/R^{\Xi^-\pi^+\bar{K}^0}$</th>
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</thead>
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<tr>
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<td>2.6</td>
<td>3.0</td>
<td>3.3</td>
<td>3.3</td>
<td>7.3</td>
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<td>1.0</td>
<td>1.0</td>
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</tr>
<tr>
<td>$\Omega(2012)$ resonance parameters</td>
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<td>33.5</td>
<td>41.3</td>
<td>27.8</td>
<td>10.7</td>
<td>41.3</td>
<td>6.1</td>
</tr>
<tr>
<td>$\Xi^0$ mass</td>
<td>...</td>
<td>...</td>
<td>17.4</td>
<td>3.3</td>
<td>...</td>
<td>17.4</td>
<td>4.5</td>
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<tr>
<td>Background parameter</td>
<td>7.9</td>
<td>2.34</td>
<td>30.0</td>
<td>17.2</td>
<td>7.9</td>
<td>30.0</td>
<td>18.1</td>
</tr>
<tr>
<td>Sum in quadrature</td>
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<td>41.0</td>
<td>54.0</td>
<td>33.0</td>
<td>13.7</td>
<td>54.1</td>
<td>21.0</td>
</tr>
</tbody>
</table>

H. 90% C.L. upper limits

In the absence of any significant observed signals, upper limits at 90% C.L. on the $R^{\Xi^-\pi^+\bar{K}^-}/R^{\Xi^-\pi^+\bar{K}^0}$, $R^{\Xi^-\pi^-\bar{K}^-}/R^{\Xi^-\pi^-\bar{K}^0}$, $R^{\Xi^-\pi^+\bar{K}^-}/R^{\Xi^-\pi^+\bar{K}^0}$, $R^{\Xi^-\pi^-\bar{K}^-}/R^{\Xi^-\pi^-\bar{K}^0}$, and $R^{\Xi^-\pi^+\bar{K}^-}/R^{\Xi^-\pi^+\bar{K}^0}$ modes are determined by solving the equation...
\[ \int_0^{1.1} \mathcal{F}_{\text{likelihood}}(t)dt / \int_0^{\infty} \mathcal{F}_{\text{likelihood}}(t)dt = 90\%, \]  
(10)

where \( t \) is the assumed ratio of branching fractions, and \( \mathcal{F}_{\text{likelihood}}(t) \) is the corresponding maximized likelihood of the data. To take into account systematic uncertainties, the likelihood is convolved with a Gaussian function whose width equals the corresponding total systematic uncertainty. Finally, we obtain

\[ R_{\Xi K}^{\Xi^{-}\pi K} = \frac{B(\Omega(2012) \rightarrow \Xi(1530)^0 \rightarrow \Xi^{-}\pi^+ K^-)}{B(\Omega(2012) \rightarrow \Xi^- K^0)} < 9.3\%, \]  
(11)

\[ R_{\Xi K}^{\Xi^{-}\pi K^0} = \frac{B(\Omega(2012) \rightarrow \Xi(1530)^- \rightarrow \Xi^{-}\pi^0 K^0)}{B(\Omega(2012) \rightarrow \Xi^- K^0)} < 81.1\%, \]  
(12)

\[ R_{\Xi K}^{\Xi^{-}\pi K^{-}} = \frac{B(\Omega(2012) \rightarrow \Xi(1530)^0 \rightarrow \Xi^{-}\pi^- K^{-})}{B(\Omega(2012) \rightarrow \Xi^- K^{-})} < 21.3\%, \]  
(13)

\[ R_{\Xi K}^{\Xi^{-}\pi K^{-}} = \frac{B(\Omega(2012) \rightarrow \Xi(1530)^0 \rightarrow \Xi^{-}\pi^0 K^-)}{B(\Omega(2012) \rightarrow \Xi^- K^-)} < 30.4\%, \]  
(14)

\[ R_{\Xi K}^{\Xi^{-}\pi K^{-}} = \frac{B(\Omega(2012) \rightarrow \Xi(1530)^0 \rightarrow \Xi^{-}\pi^- K^{-})}{B(\Omega(2012) \rightarrow \Xi^- K^-)} < 7.8\%, \]  
(15)

\[ R_{\Xi K}^{\Xi^{-}\pi K^0} = \frac{B(\Omega(2012) \rightarrow \Xi(1530)^- \rightarrow \Xi^{-}\pi^0 K^0)}{B(\Omega(2012) \rightarrow \Xi^- K^0)} < 25.6\%, \]  
(16)

and

\[ R_{\Xi K}^{\Xi^{-}\pi K^0} = \frac{B(\Omega(2012) \rightarrow \Xi(1530)^0 \rightarrow \Xi^- K^0)}{B(\Omega(2012) \rightarrow \Xi^- K^0)} < 11.9\% \]  
(17)

at 90% C.L.

### IV. RESULTS AND DISCUSSION

In summary, using the data samples of 5.7 fb\(^{-1}\) \( \Upsilon(1S) \), 24.9 fb\(^{-1}\) \( \Upsilon(2S) \), and 2.9 fb\(^{-1}\) \( \Upsilon(3S) \) collected by the Belle detector, we have searched for the three-body \( K\pi\Xi \) decay of \( \Omega(2012) \) for the first time. No significant signals are observed, and we determine upper limits at 90% C.L. on the ratios of \( R_{\Xi K}^{\Xi^{-}\pi K^0}, R_{\Xi K}^{\Xi^{-}\pi K^{-}}, R_{\Xi K}^{\Xi^{-}\pi K}, R_{\Xi K}^{\Xi^{-}\pi K^0}, R_{\Xi K}^{\Xi^{-}\pi K}, R_{\Xi K}^{\Xi^{-}\pi K}, R_{\Xi K}^{\Xi^{-}\pi K}, R_{\Xi K}^{\Xi^{-}\pi K}, \) and \( R_{\Xi K}^{\Xi^{-}\pi K} \) to be 9.3%, 81.1%, 21.3%, 30.4%, 7.8%, 25.6%, and 11.9%, respectively. Our result strongly favors the molecular interpretation proposed by Ref. [7], and is in tension with the predictions of Refs. [8–11], also based on molecular interpretations.

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[20] In common high energy physics usage, this Bayesian interval has been reported as “confidence interval,” which is a frequentist-statistics term.