

Measurements of branching fractions and asymmetry parameters of $\Xi_c^0 \rightarrow \Lambda \bar{K}^{*0}$, $\Xi_c^0 \rightarrow \Sigma^0 \bar{K}^{*0}$, and $\Xi_c^0 \rightarrow \Sigma^+ K^{*-}$ decays at Belle



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ABSTRACT: Using a data sample of 980 fb^{-1} collected with the Belle detector at the KEKB asymmetric-energy e^+e^- collider, we study the processes of $\Xi_c^0 \rightarrow \Lambda \bar{K}^{*0}$, $\Xi_c^0 \rightarrow \Sigma^0 \bar{K}^{*0}$, and $\Xi_c^0 \rightarrow \Sigma^+ K^{*-}$ for the first time. The relative branching ratios to the normalization mode of $\Xi_c^0 \rightarrow \Xi^- \pi^+$ are measured to be

$$\begin{aligned} \mathcal{B}(\Xi_c^0 \rightarrow \Lambda \bar{K}^{*0}) / \mathcal{B}(\Xi_c^0 \rightarrow \Xi^- \pi^+) &= 0.18 \pm 0.02(\text{stat.}) \pm 0.01(\text{syst.}), \\ \mathcal{B}(\Xi_c^0 \rightarrow \Sigma^0 \bar{K}^{*0}) / \mathcal{B}(\Xi_c^0 \rightarrow \Xi^- \pi^+) &= 0.69 \pm 0.03(\text{stat.}) \pm 0.03(\text{syst.}), \\ \mathcal{B}(\Xi_c^0 \rightarrow \Sigma^+ K^{*-}) / \mathcal{B}(\Xi_c^0 \rightarrow \Xi^- \pi^+) &= 0.34 \pm 0.06(\text{stat.}) \pm 0.02(\text{syst.}), \end{aligned}$$

where the uncertainties are statistical and systematic, respectively. We obtain

$$\begin{aligned} \mathcal{B}(\Xi_c^0 \rightarrow \Lambda \bar{K}^{*0}) &= (3.3 \pm 0.3(\text{stat.}) \pm 0.2(\text{syst.}) \pm 1.0(\text{ref.})) \times 10^{-3}, \\ \mathcal{B}(\Xi_c^0 \rightarrow \Sigma^0 \bar{K}^{*0}) &= (12.4 \pm 0.5(\text{stat.}) \pm 0.5(\text{syst.}) \pm 3.6(\text{ref.})) \times 10^{-3}, \\ \mathcal{B}(\Xi_c^0 \rightarrow \Sigma^+ K^{*-}) &= (6.1 \pm 1.0(\text{stat.}) \pm 0.4(\text{syst.}) \pm 1.8(\text{ref.})) \times 10^{-3}, \end{aligned}$$

where the uncertainties are statistical, systematic, and from $\mathcal{B}(\Xi_c^0 \rightarrow \Xi^- \pi^+)$, respectively. The asymmetry parameters $\alpha(\Xi_c^0 \rightarrow \Lambda \bar{K}^{*0})$ and $\alpha(\Xi_c^0 \rightarrow \Sigma^+ K^{*-})$ are $0.15 \pm 0.22(\text{stat.}) \pm 0.04(\text{syst.})$ and $-0.52 \pm 0.30(\text{stat.}) \pm 0.02(\text{syst.})$, respectively, where the uncertainties are statistical followed by systematic.

KEYWORDS: Branching fraction, e^+e^- Experiments, Charm physics

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1 Introduction

In comparison with the lowest-lying charmed baryon state, the Λ_c^+ , our knowledge of the Ξ_c states is still limited [1]. Recently, there have been many measurements of the lifetime and decay modes of the Ξ_c made by several experiments. The lifetimes of the Ξ_c^0 and Ξ_c^+ are (154.5 ± 2.5) fs and (456.8 ± 5.5) fs [2], respectively. The absolute branching fraction of $\Xi_c^0 \rightarrow \Xi^- \pi^+$ has been measured to be $(1.80 \pm 0.52)\%$, so that now the branching fractions of other channels can be determined from ratios of branching fractions [3]. The first branching fraction of the decay of the Ξ_c^0 to a charmed baryon has been measured to be $\mathcal{B}(\Xi_c^0 \rightarrow \Lambda_c^+ \pi^-) = (0.55 \pm 0.02 \pm 0.18)\%$ [4]. The branching fraction ratios of resonant and non-resonant decays of $\Xi_c^0 \rightarrow \Xi^0 K^+ K^-$ with respect to $\Xi_c^0 \rightarrow \Xi^- \pi^+$ are $0.036 \pm 0.004 \pm 0.002$ and $0.039 \pm 0.004 \pm 0.002$ [5], respectively. The branching fractions of semileptonic decays $\Xi_c^0 \rightarrow \Xi^- e^+ \nu_e$ and $\Xi_c^0 \rightarrow \Xi^- \mu^+ \nu_\mu$ have been measured, with much improved precision than hitherto, to be $(1.72 \pm 0.10 \pm 0.12 \pm 0.50)\%$ and $(1.71 \pm 0.17 \pm 0.13 \pm 0.50)\%$ [6], respectively, where the first, second, and third uncertainties are statistical, systematic, and from $\mathcal{B}(\Xi_c^0 \rightarrow \Xi^- \pi^+)$ [3]. The decay asymmetry parameter $\alpha(\Xi_c^0 \rightarrow \Xi^- \pi^+)$ has been measured to be $0.59 \pm 0.03 \pm 0.02$ [6].

Unlike semileptonic decays of Ξ_c^0 , which proceed via weak decay processes mediated by W bosons, the non-leptonic weak decays are caused by the W -boson exchanges with QCD corrections [7]. However, as the strong coupling is large at the typical energies of charm decays, it is very difficult to make quantitative predictions of decay rates and asymmetries with QCD corrections. Theoretical calculations for the hadronic decays of the Ξ_c have been performed based on $SU(3)_F$ flavor symmetry [7–25] and dynamical models [26–34]. The two-body $\mathbf{B}_c \rightarrow \mathbf{B}_n V$ decays have been calculated, where \mathbf{B}_c and \mathbf{B}_n correspond to the antitriplet charmed and light baryons, and V stands for the vector mesons, respectively.

Channel	KK [26]	Zen [31]	HYZ [24]	GLT [7]
$\Xi_c^0 \rightarrow \Lambda \bar{K}^{*0}$	1.55	1.15	0.46 ± 0.21	1.37 ± 0.26
$\Xi_c^0 \rightarrow \Sigma^0 \bar{K}^{*0}$	0.85	0.77	0.27 ± 0.22	0.42 ± 0.23
$\Xi_c^0 \rightarrow \Sigma^+ K^{*-}$	0.54	0.37	0.93 ± 0.29	0.24 ± 0.17

Table 1. Decay branching fractions (%) of the Cabibbo-favored $\Xi_c^0 \rightarrow \mathbf{B}_n V$ decays based on the covariant quark model from KK [26], pole model from Zen [31], and $SU(3)_F$ flavor symmetry from HYZ [24] and GLT [7].

Channel	KK [26]	Zen [31]	GLT [7]
$\Xi_c^0 \rightarrow \Lambda \bar{K}^{*0}$	+0.58	+0.49	-0.67 ± 0.24
$\Xi_c^0 \rightarrow \Sigma^0 \bar{K}^{*0}$	-0.87	+0.25	-0.42 ± 0.62
$\Xi_c^0 \rightarrow \Sigma^+ K^{*-}$	-0.60	+0.51	$-0.76_{-0.24}^{+0.64}$

Table 2. The asymmetry parameters for the Cabibbo-favored $\Xi_c^0 \rightarrow \mathbf{B}_n V$ decays based on the covariant quark model from KK [26], pole model from Zen [31], and $SU(3)_F$ flavor symmetry from GLT [7].

However, the different models give widely varying predictions. For $\Xi_c^0 \rightarrow \Lambda \bar{K}^{*0}$, $\Xi_c^0 \rightarrow \Sigma^0 \bar{K}^{*0}$, and $\Xi_c^0 \rightarrow \Sigma^+ K^{*-}$, the predicted branching fractions cover wide ranges of (0.46 – 1.55)%, (0.27 – 0.85)%, and (0.24 – 0.93)% [7, 24, 26, 31], as listed in table 1 for three different models.

For the channels $\Xi_c^0 \rightarrow \Lambda \bar{K}^{*0}$, $\Xi_c^0 \rightarrow \Sigma^0 \bar{K}^{*0}$, and $\Xi_c^0 \rightarrow \Sigma^+ K^{*-}$, parity violation is manifested by polarization of the hyperon (Λ , Σ^0 , or Σ^+), and is quantified by the decay asymmetry parameter, α . Because the hyperon decay also violates parity, the product of decay asymmetry parameters of Ξ_c^0 decay and hyperon decay can be measured by its decay angular distribution. Note that the asymmetry parameter of $\Xi_c^0 \rightarrow \Sigma^0 \bar{K}^{*0}$ can not be measured since the value of $\alpha(\Sigma^0 \rightarrow \gamma \Lambda)$ should be zero for an electromagnetic decay of $\Sigma^0 \rightarrow \gamma \Lambda$. We measure the product value of $\alpha(\Xi_c^0 \rightarrow \Sigma^0 \bar{K}^{*0})\alpha(\Sigma^0 \rightarrow \gamma \Lambda)$ just to validate no bias in the measurement. Different models produce widely different predictions for α [7, 26, 31], which are listed in table 2.

In this article, we use the entire data sample of 980 fb^{-1} integrated luminosity collected at Belle to perform the first measurements of the branching fractions and asymmetry parameters for the decays $\Xi_c^0 \rightarrow \Lambda \bar{K}^{*0}$, $\Xi_c^0 \rightarrow \Sigma^0 \bar{K}^{*0}$, and $\Xi_c^0 \rightarrow \Sigma^+ K^{*-}$. The Λ , \bar{K}^{*0} , Σ^0 , Σ^+ , and K^{*-} are reconstructed by $p\pi^-$, $K^-\pi^+$, $\Lambda\gamma$, $p\pi^0$, and $K_S^0\pi^-$ final states, respectively. Throughout this analysis, for any given mode, the corresponding charge-conjugate mode is implied.

2 Data sample and the Belle detector

This measurement is based on data recorded at or near the $\Upsilon(1S)$, $\Upsilon(2S)$, $\Upsilon(3S)$, $\Upsilon(4S)$, and $\Upsilon(5S)$ resonances by the Belle detector [35, 36] at the KEKB asymmetric-energy e^+e^- collider [37, 38]. The total data sample corresponds to an integrated luminosity

of 980 fb^{-1} [36]. The Belle detector is a large solid-angle magnetic spectrometer that consists of a silicon vertex detector, a 50-layer central drift chamber (CDC), an array of aerogel threshold Cherenkov counters (ACC), a barrel-like arrangement of time-of-flight scintillation counters (TOF), and an electromagnetic calorimeter comprised of CsI(Tl) crystals (ECL) located inside a superconducting solenoid coil that provides a 1.5T magnetic field. An iron flux-return yoke instrumented with resistive plate chambers located outside the coil is used to detect K_L^0 mesons and identify muons. A detailed description of the Belle detector can be found in refs. [35, 36].

Samples of simulated signal events are generated using EVTGEN [39] to optimize the signal selection criteria and calculate the signal reconstruction efficiency; $e^+e^- \rightarrow c\bar{c}$ events are simulated using PYTHIA [40], and $\Xi_c^0 \rightarrow \Lambda \bar{K}^{*0}/\Sigma^0 \bar{K}^{*0}/\Sigma^+ K^{*-}$ decays are generated with a phase space model. The effect of final-state radiation is taken into account in the simulation using the PHOTOS [41] package. The simulated events are processed with a detector simulation based on GEANT3 [42]. Generic simulated samples, i.e. $B = B^+, B^0$, or $B_s^{(*)}$ decays and $e^+e^- \rightarrow q\bar{q}$ ($q = u, d, s, c$) at $\sqrt{s} = 10.52, 10.58, \text{ and } 10.867 \text{ GeV}$, normalized to the same integrated luminosity as real data, are used to check peaking backgrounds and to perform input/output checks.

3 Selection criteria

Except for the charged tracks from the relatively long-lived $\Lambda \rightarrow p\pi^-$ and $K_S^0 \rightarrow \pi^+\pi^-$ decays, impact parameters with respect to the interaction point (IP) are required to be less than 1 cm and 4 cm perpendicular to, and along the beam direction, respectively. For the particle identification (PID) of a well-reconstructed charged track, information from different detector subsystems, including specific ionization in the CDC, time measurement in the TOF, and the response of the ACC, is combined to form a likelihood \mathcal{L}_i [43] for particle species i . Tracks with $R_K = \mathcal{L}_K/(\mathcal{L}_K + \mathcal{L}_\pi) < 0.4$ are identified as pions with an efficiency of 95%, while 6% of kaons are misidentified as pions; tracks with $R_K > 0.6$ are identified as kaons with an efficiency of 96%, while 7% of pions are misidentified as kaons. For proton identification, a track with $R_p^\pi = \mathcal{L}_p/(\mathcal{L}_p + \mathcal{L}_\pi) > 0.6$ and $R_p^K = \mathcal{L}_p/(\mathcal{L}_p + \mathcal{L}_K) > 0.6$ is identified as a proton with an efficiency of about 98%; less than 1% of the pions/kaons are misidentified as protons.

Using a multivariate analysis with a neural network [44] based on two sets of input variables [45], a K_S^0 candidate is reconstructed from a pair of oppositely charged tracks that are treated as pions. The invariant mass of the K_S^0 candidates is required to be within $10 \text{ MeV}/c^2$ of the corresponding nominal mass [1]. Candidate Λ baryons are reconstructed in the decay $\Lambda \rightarrow p\pi^-$ and selected if $|M(p\pi^-) - m_\Lambda| < 3 \text{ MeV}/c^2$ ($\sim 2.5\sigma$), where σ denotes the mass resolution. Hereinafter, M represents a measured invariant mass and m_i denotes the nominal mass of particle i [1].

An ECL cluster is treated as a photon candidate if it is isolated from the projected path of charged tracks in the CDC, and its energy in the laboratory frame is greater than 50 MeV. To reduce photon candidates originating from neutral hadrons, we reject a

photon candidate if the ratio of energy deposited in the central 3×3 square of cells to that deposited in the enclosing 5×5 square of cells in its ECL cluster is less than 0.85.

For the photons from $\Sigma^0 \rightarrow \gamma\Lambda$ decays, we require $E_\gamma > 0.1 \text{ GeV}$ in the laboratory frame to suppress backgrounds. Photon pairs are kept as π^0 candidates if their invariant mass lies in the range $120 \text{ MeV}/c^2 < M(\gamma\gamma) < 150 \text{ MeV}/c^2$ ($\pm 3\sigma$ about the nominal mass of π^0). In the $\Sigma^+ \rightarrow p\pi^0$ reconstruction, combinations of π^0 candidates and protons are made using those protons with a significantly large ($> 1 \text{ mm}$) impact parameter with respect to the IP. The flight directions of Σ^+ candidates, which are reconstructed from their fitted production and decay vertices, are required to be consistent with their momentum directions [46].

The $\Lambda\bar{K}^{*0}$, $\Sigma^0\bar{K}^{*0}$, or Σ^+K^{*-} candidates are combined to form a Ξ_c^0 with its daughter tracks fitted to a common vertex. To reduce combinatorial backgrounds, especially from B -meson decays, the scaled momentum $x_p = p^*/p_{\text{max}}$ is required to be greater than 0.5. Here, p^* is the momentum of Ξ_c^0 in the center-of-mass (C.M.) frame, and $p_{\text{max}} = \sqrt{E_{\text{beam}}^2 - M_{\Xi_c^0}^2}c^4/c$ is the maximum momentum, where E_{beam} is the beam energy.

We veto $\Sigma(1385)$ intermediate backgrounds by requiring $M(\Lambda\pi^+) > 1.42 \text{ GeV}/c^2$ and $M(\Sigma^0\pi^+/\Sigma^+\pi^-) > 1.45 \text{ GeV}/c^2$ in the substructures of $\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0}$ and $\Xi_c^0 \rightarrow \Sigma^0\bar{K}^{*0}/\Sigma^+K^{*-}$ candidates, respectively. These requirements can also suppress backgrounds from D^{*+} decays, for which the momentum of pions from D^{*+} is low. No peaking backgrounds are found in $M(\Lambda K^-)$, $M(\Sigma^0 K^-)$, and $M(\Sigma^+ K_S^0)$ distributions from generic simulated samples [47].

To determine the absolute branching fractions of $\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0}$, $\Xi_c^0 \rightarrow \Sigma^0\bar{K}^{*0}$, and $\Xi_c^0 \rightarrow \Sigma^+K^{*-}$, the reference mode of $\Xi_c^0 \rightarrow \Xi^-\pi^+$ is utilized. Selections of candidates in $\Xi_c^0 \rightarrow \Xi^-\pi^+$ use well-reconstructed tracks, PID, and the vertex fitting technique in a way similar to the methods in ref. [6]. Candidate Λ baryons are reconstructed as above. We define the Ξ^- signal region as $|M(\Lambda\pi^-) - m_{\Xi^-}| < 6.5 \text{ MeV}/c^2$ ($\sim 3\sigma$). To suppress the combinational background, we require the flight directions of Λ and Ξ^- candidates, which are reconstructed from their fitted production and decay vertices, to be within five degrees of their momentum directions. We also require the scaled momentum $x_p > 0.5$.

4 Branching fractions of $\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0}$, $\Xi_c^0 \rightarrow \Sigma^0\bar{K}^{*0}$, and $\Xi_c^0 \rightarrow \Sigma^+K^{*-}$

After applying the above requirements, the invariant mass distributions of $p\pi^-$, $\Lambda\gamma$, and $p\pi^0$ from data samples are shown in figure 1. The Λ , Σ^0 , and Σ^+ signals are clear. We define Λ , Σ^0 , and Σ^+ signal regions as $|M(p\pi^-) - m_\Lambda| < 3 \text{ MeV}/c^2$ ($\sim 2.5\sigma$), $|M(\Lambda\gamma) - m_{\Sigma^0}| < 12 \text{ MeV}/c^2$ ($\sim 2.5\sigma$), and $|M(p\pi^0) - m_{\Sigma^+}| < 12 \text{ MeV}/c^2$ ($\sim 2.5\sigma$). The mass sideband regions, which will be taken to study backgrounds to our signals, are twice as wide as the signal regions, as shown by the blue dashed lines in figure 1.

In the $\Lambda\bar{K}^{*0}$, $\Sigma^0\bar{K}^{*0}$, and Σ^+K^{*-} invariant mass spectra, no peaking background is found from combinatorial backgrounds of Λ , Σ^0 , and Σ^+ candidates, but we found a fraction of Ξ_c^0 signal events from the events outside of the \bar{K}^{*0} and K^{*-} signal regions, especially in the higher side of $M(K^-\pi^+)$ and $M(K_S^0\pi^-)$, as shown in figure 2.

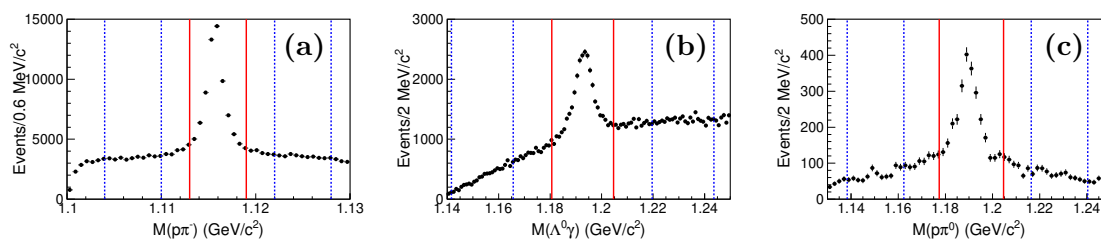


Figure 1. The invariant mass distributions for Λ , Σ^0 , and Σ^+ candidates from data samples. The red solid lines show the required signal regions, and the blue dashed lines show the defined mass sidebands.

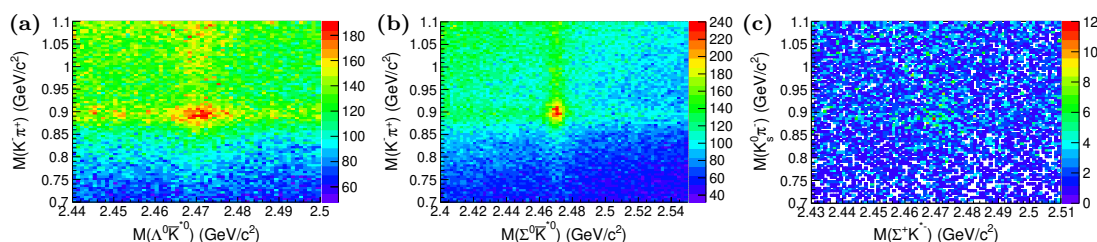


Figure 2. The scatter plots of (a) $M(K^-\pi^+)$ versus $M(\Lambda\bar{K}^{*0})$, (b) $M(K^-\pi^+)$ versus $M(\Sigma^0\bar{K}^{*0})$, and (c) $M(K_S^0\pi^-)$ versus $M(\Sigma^+K^{*-})$ distributions from data samples.

To extract the Ξ_c^0 signal yields from two-body $\Lambda\bar{K}^{*0}/\Sigma^0\bar{K}^{*0}/\Sigma^+K^{*-}$ decay, we perform a two-dimensional (2D) binned maximum-likelihood fit to $M(K^-\pi^+/K^-\pi^+/K_S^0\pi^-)$ and $M(\Lambda\bar{K}^{*0}/\Sigma^0\bar{K}^{*0}/\Sigma^+K^{*-})$ distributions assuming there is no interference between the $\bar{K}^{*0}/\bar{K}^{*0}/K^{*-}$ signals and backgrounds. The 2D fitting function $f(M_1, M_2)$ is expressed as

$$f(M_1, M_2) = N^{\text{sig}} s_1(M_1) s_2(M_2) + N_{\text{sb}}^{\text{bg}} s_1(M_1) b_2(M_2) + N_{\text{bs}}^{\text{bg}} b_1(M_1) s_2(M_2) + N_{\text{bb}}^{\text{bg}} b_1(M_1) b_2(M_2), \quad (4.1)$$

where $s_1(M_1)$ and $b_1(M_1)$ are the signal and background probability density functions (PDFs) for the $M(\Lambda\bar{K}^{*0}/\Sigma^0\bar{K}^{*0}/\Sigma^+K^{*-})$ distributions, respectively, and $s_2(M_2)$ and $b_2(M_2)$ are the corresponding PDFs for the $M(K^-\pi^+/K^-\pi^+/K_S^0\pi^-)$ distributions. Here, $N_{\text{sb}}^{\text{bg}}$ and $N_{\text{bs}}^{\text{bg}}$ denote the numbers of peaking background events in $M(\Lambda\bar{K}^{*0}/\Sigma^0\bar{K}^{*0}/\Sigma^+K^{*-})$ and $M(K^-\pi^+/K^-\pi^+/K_S^0\pi^-)$ distributions, and $N_{\text{bb}}^{\text{bg}}$ is the number of combinatorial background in both Ξ_c^0 and \bar{K}^{*0}/K^{*-} candidates. The signal shapes of Ξ_c^0 and \bar{K}^{*0}/K^{*-} ($s_1(M_1)$ and $s_2(M_2)$) are described by Breit-Wigner (BW) functions convolved with Gaussian functions, and second- or third-order polynomial functions represent the backgrounds ($b_1(M_1)$ and $b_2(M_2)$). The values of signal PDF parameters are fixed to those obtained from the fits to the corresponding simulated signal distributions. The values of the background shape parameters are free. The fitted results are shown in figure 3. The fitted Ξ_c^0 signal yields in $\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0}$, $\Xi_c^0 \rightarrow \Sigma^0\bar{K}^{*0}$, and $\Xi_c^0 \rightarrow \Sigma^+K^{*-}$ are 3974 ± 367 , 6260 ± 254 , and 373 ± 61 events, respectively.

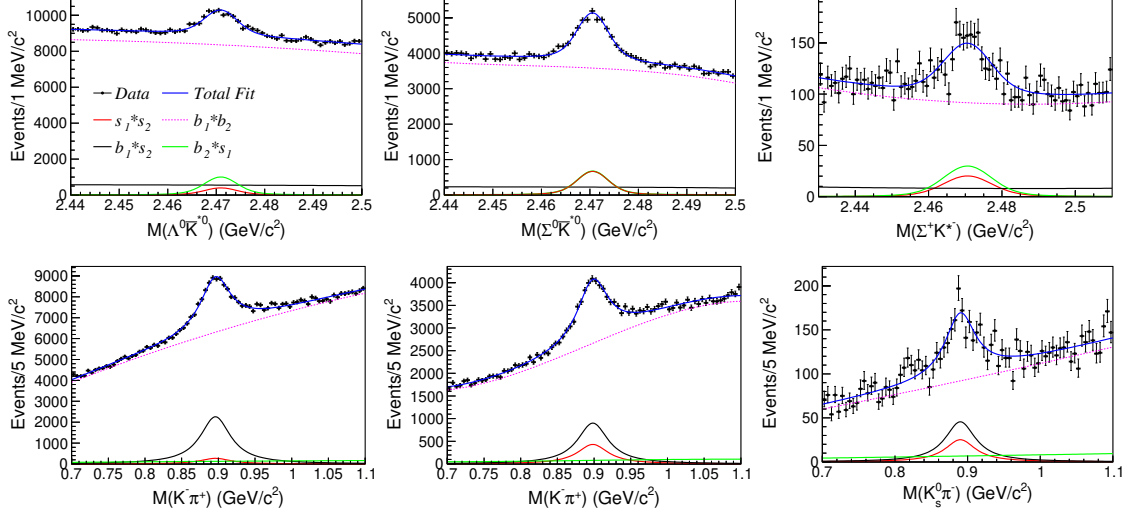


Figure 3. 2D binned maximum-likelihood fits to $M(K^-\pi^+)$ and $M(\Lambda\bar{K}^{*0})$ (left), $M(K^-\pi^+)$ and $M(\Sigma^0\bar{K}^{*0})$ (middle), and $M(K_S^0\pi^-)$ and $M(\Sigma^+K^{*-})$ (right) distributions from data samples. All components are indicated in the labels and described in the text.

We calculate the branching fraction ratios according to the equations

$$\frac{\mathcal{B}(\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0})}{\mathcal{B}(\Xi_c^0 \rightarrow \Xi^-\pi^+)} = \frac{N_{\Lambda\bar{K}^{*0}}^{\text{sig}} \cdot \varepsilon_{\Xi^-\pi^+} \cdot \mathcal{B}(\Xi^- \rightarrow \Lambda\pi^-)}{N_{\Xi^-\pi^+}^{\text{sig}} \cdot \varepsilon_{\Lambda\bar{K}^{*0}} \cdot \mathcal{B}(\bar{K}^{*0} \rightarrow K^-\pi^+)}, \quad (4.2)$$

$$\frac{\mathcal{B}(\Xi_c^0 \rightarrow \Sigma^0\bar{K}^{*0})}{\mathcal{B}(\Xi_c^0 \rightarrow \Xi^-\pi^+)} = \frac{N_{\Sigma^0\bar{K}^{*0}}^{\text{sig}} \cdot \varepsilon_{\Xi^-\pi^+} \cdot \mathcal{B}(\Xi^- \rightarrow \Lambda\pi^-)}{N_{\Xi^-\pi^+}^{\text{sig}} \cdot \varepsilon_{\Sigma^0\bar{K}^{*0}} \cdot \mathcal{B}(\Sigma^0 \rightarrow \gamma\Lambda) \cdot \mathcal{B}(\bar{K}^{*0} \rightarrow K^-\pi^+)}, \quad (4.3)$$

and

$$\frac{\mathcal{B}(\Xi_c^0 \rightarrow \Sigma^+K^{*-})}{\mathcal{B}(\Xi_c^0 \rightarrow \Xi^-\pi^+)} = \frac{N_{\Sigma^+K^{*-}}^{\text{sig}} \cdot \varepsilon_{\Xi^-\pi^+}}{N_{\Xi^-\pi^+}^{\text{sig}} \cdot \varepsilon_{\Sigma^+K^{*-}}} \times \frac{\mathcal{B}(\Lambda \rightarrow p\pi^-)}{\mathcal{B}(\Sigma^+ \rightarrow p\pi^0) \cdot \mathcal{B}(\pi^0 \rightarrow \gamma\gamma) \cdot \mathcal{B}(K^{*-} \rightarrow K_S^0\pi^-) \cdot \mathcal{B}(K_S^0 \rightarrow \pi^+\pi^-)}. \quad (4.4)$$

Here, $N_{\Lambda\bar{K}^{*0}}^{\text{sig}}$, $N_{\Sigma^0\bar{K}^{*0}}^{\text{sig}}$, $N_{\Sigma^+K^{*-}}^{\text{sig}}$, and $N_{\Xi^-\pi^+}^{\text{sig}}$ are the signal yields, and $\varepsilon_{\Lambda\bar{K}^{*0}} = 18.9\%$, $\varepsilon_{\Sigma^0\bar{K}^{*0}} = 8.0\%$, $\varepsilon_{\Sigma^+K^{*-}} = 3.5\%$, and $\varepsilon_{\Xi^-\pi^+} = 27.9\%$ are reconstruction efficiencies found from the signal simulations. In eqs. (4.2) and (4.3), the same intermediate branching fraction of $\mathcal{B}(\Lambda \rightarrow p\pi^-)$ is canceled since the Λ candidates are all reconstructed by $p\pi^-$. Branching fractions $\mathcal{B}(\Xi^- \rightarrow \Lambda\pi^-) = (99.887 \pm 0.035)\%$, $\mathcal{B}(\Sigma^0 \rightarrow \gamma\Lambda) = 100\%$, $\mathcal{B}(\bar{K}^{*0} \rightarrow K^-\pi^+) = 66.7\%$, $\mathcal{B}(\Lambda \rightarrow p\pi^-) = (63.9 \pm 0.5)\%$, $\mathcal{B}(\Sigma^+ \rightarrow p\pi^0) = (51.57 \pm 0.30)\%$, $\mathcal{B}(\pi^0 \rightarrow \gamma\gamma) = (98.823 \pm 0.034)\%$, $\mathcal{B}(K^{*-} \rightarrow K_S^0\pi^-) = 33.34\%$, and $\mathcal{B}(K_S^0 \rightarrow \pi^+\pi^-) = (69.2 \pm 0.05)\%$ are from PDG [1] directly or calculated based on the isospin symmetry. The branching fraction of $\Xi_c^0 \rightarrow \Xi^-\pi^+$ is $(1.80 \pm 0.52)\%$ [3]. Using the values above, we measure relative branching ratios to the normalization mode of $\Xi_c^0 \rightarrow \Xi^-\pi^+$ and the branching fractions of $\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0}$, $\Xi_c^0 \rightarrow \Sigma^0\bar{K}^{*0}$, and $\Xi_c^0 \rightarrow \Sigma^+K^{*-}$, which are summarized in table 3.

$\mathcal{B}(\Xi_c^0 \rightarrow \Lambda \bar{K}^{*0})/\mathcal{B}(\Xi_c^0 \rightarrow \Xi^- \pi^+)$	$0.18 \pm 0.02(\text{stat.}) \pm 0.01(\text{syst.})$
$\mathcal{B}(\Xi_c^0 \rightarrow \Lambda \bar{K}^{*0})$	$(3.3 \pm 0.3(\text{stat.}) \pm 0.2(\text{syst.}) \pm 1.0(\text{ref.})) \times 10^{-3}$
$\mathcal{B}(\Xi_c^0 \rightarrow \Sigma^0 \bar{K}^{*0})/\mathcal{B}(\Xi_c^0 \rightarrow \Xi^- \pi^+)$	$0.69 \pm 0.03(\text{stat.}) \pm 0.03(\text{syst.})$
$\mathcal{B}(\Xi_c^0 \rightarrow \Sigma^0 \bar{K}^{*0})$	$(12.4 \pm 0.5(\text{stat.}) \pm 0.5(\text{syst.}) \pm 3.6(\text{ref.})) \times 10^{-3}$
$\mathcal{B}(\Xi_c^0 \rightarrow \Sigma^+ K^{*-})/\mathcal{B}(\Xi_c^0 \rightarrow \Xi^- \pi^+)$	$0.34 \pm 0.06(\text{stat.}) \pm 0.02(\text{syst.})$
$\mathcal{B}(\Xi_c^0 \rightarrow \Sigma^+ K^{*-})$	$(6.1 \pm 1.0(\text{stat.}) \pm 0.4(\text{syst.}) \pm 1.8(\text{ref.})) \times 10^{-3}$

Table 3. The branching fractions and ratios, where the uncertainties are statistical, systematic, and from $\mathcal{B}(\Xi_c^0 \rightarrow \Xi^- \pi^+)$ [3].

5 Asymmetry parameter extraction

For $\Xi_c^0 \rightarrow \Lambda \bar{K}^{*0}$, the differential decay rate [7] is given by:

$$\frac{dN}{d\cos\theta_\Lambda} \propto 1 + \alpha(\Xi_c^0 \rightarrow \Lambda \bar{K}^{*0})\alpha(\Lambda \rightarrow p\pi^-)\cos\theta_\Lambda, \quad (5.1)$$

where $\alpha(\Xi_c^0 \rightarrow \Lambda \bar{K}^{*0})$ and $\alpha(\Lambda \rightarrow p\pi^-)$ are the asymmetry parameters of $\Xi_c^0 \rightarrow \Lambda \bar{K}^{*0}$ and $\Lambda \rightarrow p\pi^-$, and θ_Λ is the angle between the proton momentum vector and the opposite of Ξ_c^0 momentum vector in the Λ rest frame.

For $\Xi_c^0 \rightarrow \Sigma^0 \bar{K}^{*0}$, the differential decay rate [7] is shown as:

$$\frac{dN}{d\cos\theta_{\Sigma^0}} \propto 1 + \alpha(\Xi_c^0 \rightarrow \Sigma^0 \bar{K}^{*0})\alpha(\Sigma^0 \rightarrow \gamma\Lambda)\cos\theta_{\Sigma^0}, \quad (5.2)$$

where $\alpha(\Xi_c^0 \rightarrow \Sigma^0 \bar{K}^{*0})$ and $\alpha(\Sigma^0 \rightarrow \gamma\Lambda)$ are the asymmetry parameters for $\Xi_c^0 \rightarrow \Sigma^0 \bar{K}^{*0}$ and $\Sigma^0 \rightarrow \gamma\Lambda$, and θ_{Σ^0} is the angle between the Λ momentum vector and the opposite of Ξ_c^0 momentum vector in the Σ^0 rest frame. The value of $\alpha(\Sigma^0 \rightarrow \gamma\Lambda)$ should be zero due to the conservation of parity for an electromagnetic decay. Thus, the distribution of $\cos\theta_{\Sigma^0}$ is expected to be flat.

For $\Xi_c^0 \rightarrow \Sigma^+ K^{*-}$, the differential decay rate [7] can be described with:

$$\frac{dN}{d\cos\theta_{\Sigma^+}} \propto 1 + \alpha(\Xi_c^0 \rightarrow \Sigma^+ K^{*-})\alpha(\Sigma^+ \rightarrow p\pi^0)\cos\theta_{\Sigma^+}, \quad (5.3)$$

where $\alpha(\Xi_c^0 \rightarrow \Sigma^+ K^{*-})$ and $\alpha(\Sigma^+ \rightarrow p\pi^0)$ are asymmetry parameters for $\Xi_c^0 \rightarrow \Sigma^+ K^{*-}$ and $\Sigma^+ \rightarrow p\pi^0$, and θ_{Σ^+} is the angle between the p momentum vector and the opposite of Ξ_c^0 momentum vector in the Σ^+ rest frame.

We determine the asymmetry parameters by fitting the decay angular distributions of $\cos\theta_\Lambda$, $\cos\theta_{\Sigma^0}$, and $\cos\theta_{\Sigma^+}$ with eqs. (5.1)–(5.3). For $\Xi_c^0 \rightarrow \Lambda \bar{K}^{*0}$ and $\Xi_c^0 \rightarrow \Sigma^0 \bar{K}^{*0}$, we do 2D fits as above to data in 8 $\cos\theta_\Lambda$ and $\cos\theta_{\Sigma^0}$ bins. For $\Xi_c^0 \rightarrow \Sigma^+ K^{*-}$, we do 2D fits as above to data in 5 $\cos\theta_{\Sigma^+}$ bins. The signal yields in data are summarized in tables 4 and 5. We then make the efficiency-corrected $\cos\theta_\Lambda$, $\cos\theta_{\Sigma^0}$, and $\cos\theta_{\Sigma^+}$ distributions of data, as shown in figure 4. The fitted results are shown by the lines in figure 4. The returned values of product asymmetry parameters are listed in table 6. The asymmetry parameter

$\cos\theta$	$[-1, -0.75)$	$[-0.75, -0.5)$	$[-0.5, -0.25)$	$[-0.25, 0)$	$[0, 0.25)$	$[0.25, 0.5)$	$[0.5, 0.75)$	$[0.75, 1]$
$\Lambda\bar{K}^{*0}$	$\frac{582.1\pm128.2}{0.201}$	$\frac{399.8\pm123.8}{0.196}$	$\frac{456.1\pm124.0}{0.192}$	$\frac{310.3\pm124.7}{0.187}$	$\frac{644.3\pm128.4}{0.182}$	$\frac{552.7\pm132.1}{0.182}$	$\frac{477.0\pm136.8}{0.185}$	$\frac{551.2\pm144.3}{0.191}$
$\Sigma^0\bar{K}^{*0}$	$\frac{850.0\pm85.0}{0.094}$	$\frac{905.8\pm86.0}{0.090}$	$\frac{885.4\pm87.3}{0.088}$	$\frac{837.9\pm87.4}{0.083}$	$\frac{791.6\pm89.3}{0.078}$	$\frac{782.7\pm90.5}{0.075}$	$\frac{624.8\pm92.1}{0.070}$	$\frac{581.5\pm93.3}{0.063}$

Table 4. The values of $\frac{\text{Signal yield}}{\text{Reconstruction efficiency}}$ from data samples in different $\cos\theta_\Lambda$ and $\cos\theta_{\Sigma^0}$ bins in $\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0}$ and $\Xi_c^0 \rightarrow \Sigma^0\bar{K}^{*0}$.

$\cos\theta$	$[-1, -0.6)$	$[-0.6, -0.2)$	$[-0.2, 0.2)$	$[0.2, 0.6)$	$[0.6, 1]$
Σ^+K^{*-}	$\frac{44.1\pm28.1}{0.039}$	$\frac{50.9\pm25.9}{0.033}$	$\frac{88.4\pm27.1}{0.032}$	$\frac{97.7\pm28.8}{0.033}$	$\frac{92.1\pm28.4}{0.035}$

Table 5. The values of $\frac{\text{Signal yield}}{\text{Reconstruction efficiency}}$ from data samples in different $\cos\theta_{\Sigma^+}$ bins in $\Xi_c^0 \rightarrow \Sigma^+K^{*-}$.

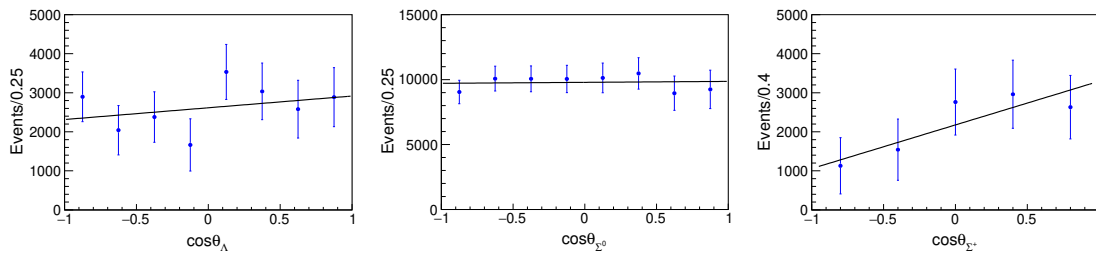


Figure 4. The $\cos\theta_\Lambda$, $\cos\theta_{\Sigma^0}$, and $\cos\theta_{\Sigma^+}$ distributions after efficiency corrections from data samples. The lines show the fitted result with functions of eq. (5.1), eq. (5.2), and eq. (5.3).

$\alpha(\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0})\alpha(\Lambda \rightarrow p\pi^-)$	$0.115 \pm 0.164(\text{stat.}) \pm 0.031(\text{syst.})$
$\alpha(\Xi_c^0 \rightarrow \Sigma^0\bar{K}^{*0})\alpha(\Sigma^0 \rightarrow \gamma\Lambda)$	$0.008 \pm 0.072(\text{stat.}) \pm 0.008(\text{syst.})$
$\alpha(\Xi_c^0 \rightarrow \Sigma^+K^{*-})\alpha(\Sigma^+ \rightarrow p\pi^0)$	$0.514 \pm 0.295(\text{stat.}) \pm 0.012(\text{syst.})$
$\alpha(\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0})$	$0.15 \pm 0.22(\text{stat.}) \pm 0.04(\text{syst.})$
$\alpha(\Xi_c^0 \rightarrow \Sigma^+K^{*-})$	$-0.52 \pm 0.30(\text{stat.}) \pm 0.02(\text{syst.})$

Table 6. The values of asymmetry parameters, where the uncertainties are statistical and systematic.

$\alpha(\Xi_c^0 \rightarrow \Sigma^0\bar{K}^{*0})\alpha(\Sigma^0 \rightarrow \gamma\Lambda)$ is $0.008 \pm 0.072(\text{stat.}) \pm 0.008(\text{syst.})$, which is consistent with zero, and compatible with parity conservation for an electromagnetic decay of $\Sigma^0 \rightarrow \gamma\Lambda$. Noting that $\alpha(\Lambda \rightarrow p\pi^-) = 0.747 \pm 0.010$ and $\alpha(\Sigma^+ \rightarrow p\pi^0) = -0.980 \pm 0.017$ from PDG [1], the $\alpha(\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0})$ and $\alpha(\Xi_c^0 \rightarrow \Sigma^+K^{*-})$ values are obtained and listed in table 6. The systematic uncertainties are discussed below.

6 Systematic uncertainties

The sources of systematic uncertainties for the branching fraction measurements include detection efficiency uncertainties, uncertainties in the branching fractions of the intermediate states, and the uncertainty associated with the fitting technique used. Note that partial uncertainties from detection efficiency sources and branching fractions can be can-

Final state	$\Lambda\bar{K}^{*0}$	$\Sigma^0\bar{K}^{*0}$	Σ^+K^{*-}
Detection efficiency	2.0	2.9	3.8
Branching fraction	—	—	1.0
Fit uncertainty	3.8	3.1	5.4
Sum in quadrature	4.3	4.2	6.7

Table 7. Relative systematic uncertainties (%) in the branching fraction measurements. The uncertainty of 28.9% on $\mathcal{B}(\Xi_c^0 \rightarrow \Xi^- \pi^+)$ [3] is treated as an independent systematic uncertainty.

celed by the reference mode of $\Xi_c^0 \rightarrow \Xi^- \pi^+$. The detection efficiency uncertainties include those from tracking efficiency (0.35%/track), particle identification efficiency (1.1%/kaon, 0.9%/pion, and 2.9%/proton), K_S^0 selection efficiency (2.23%), π^0 reconstruction efficiency (2.25%/ π^0), and photon reconstruction efficiency (2.0%/photon). The total detection efficiency uncertainties are obtained by adding all sources in quadrature.

As the partial uncertainties from branching fractions are canceled in the ratio to the reference mode, only the uncertainties of $\mathcal{B}(\Sigma^+ \rightarrow p\pi^0)$ (0.6%) and $\mathcal{B}(\Lambda^0 \rightarrow p\pi^-)$ (0.8%) [1] need to be included for $\Xi_c^0 \rightarrow \Sigma^+ K^{*-}$. In 2D fitting to $M(K^- \pi^+)$ and $M(\Lambda\bar{K}^{*0})$, $M(K^- \pi^+)$ and $M(\Sigma^0\bar{K}^{*0})$, and $M(K_S^0\pi^-)$ and $M(\Sigma^+K^{*-})$ distributions for $\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0}$, $\Xi_c^0 \rightarrow \Sigma^0\bar{K}^{*0}$, and $\Xi_c^0 \rightarrow \Sigma^+K^{*-}$, we enlarge the mass resolution by 10%, and change the fit range and background shape, then the differences of signal yields are taken as the systematic uncertainties. The total fit uncertainty is obtained by summing the uncertainties from $\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0}/\Sigma^0\bar{K}^{*0}/\Sigma^+K^{*-}$ and reference mode of $\Xi_c^0 \rightarrow \Xi^- \pi^+$ in quadrature. All the uncertainties are summarized in table 7. Finally, assuming all the sources are independent and adding them in quadrature, the total systematic uncertainties on the branching fraction measurements are calculated. The uncertainty of 28.9% on $\mathcal{B}(\Xi_c^0 \rightarrow \Xi^- \pi^+)$ [3] is treated as systematic uncertainty separately.

The sources of the systematic uncertainties in the asymmetry parameter extractions include fitting procedures, the numbers of $\cos\theta$ bins, uncertainties in $\alpha(\Lambda \rightarrow p\pi^-)$ and $\alpha(\Sigma^+ \rightarrow p\pi^0)$ values, and production polarization of Ξ_c^0 . The fitting uncertainties are estimated using simulated pseudoexperiments. We use an ensemble of simulated experiments to generate the mass spectra of Ξ_c^0 and \bar{K}^* candidates corresponding to figure 3. The number of signal events in each $\cos\theta$ bin is obtained by a 2D fit to the generated mass spectra after enlarging the mass resolution by 10%, and changing the fit range and background shape. After 10,000 simulations, distributions of $\alpha(\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0})\alpha(\Lambda \rightarrow p\pi^-)$, $\alpha(\Xi_c^0 \rightarrow \Sigma^0\bar{K}^{*0})\alpha(\Sigma^0 \rightarrow \gamma\Lambda)$, and $\alpha(\Xi_c^0 \rightarrow \Sigma^+K^{*-})\alpha(\Sigma^+ \rightarrow p\pi^0)$ are obtained by fitting the slopes of the $\cos\theta_\Lambda$, $\cos\theta_{\Sigma^0}$ and $\cos\theta_{\Sigma^+}$ distributions. The differences between the fitted peaking values of the distributions of $\alpha(\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0})\alpha(\Lambda \rightarrow p\pi^-)$, $\alpha(\Xi_c^0 \rightarrow \Sigma^0\bar{K}^{*0})\alpha(\Sigma^0 \rightarrow \gamma\Lambda)$, and $\alpha(\Xi_c^0 \rightarrow \Sigma^+K^{*-})\alpha(\Sigma^+ \rightarrow p\pi^0)$ and the nominal values are taken as fitting uncertainties.

We change the numbers of $\cos\theta$ bins from 8 and 5 to 10 and 8 for $\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0}$, $\Sigma^0\bar{K}^{*0}$ and $\Xi_c^0 \rightarrow \Sigma^+K^{*-}$, and the differences on the asymmetry parameters are taken as the related systematic uncertainty. The absolute uncertainties on $\alpha(\Lambda \rightarrow p\pi^-)$ and $\alpha(\Sigma^+ \rightarrow p\pi^0)$ values are 0.002 and 0.009 [1]. This measurement is insensitive to production polarization

Final state	$\Lambda\bar{K}^{*0}$	$\Sigma^0\bar{K}^{*0}$	Σ^+K^{*-}
Fit uncertainty	0.023	0.006	0.009
$\cos\theta$ bins	0.033	0.005	0.007
$\alpha(\Lambda \rightarrow p\pi^-)$ and $\alpha(\Sigma^+ \rightarrow p\pi^0)$ values	0.002	—	0.009
Sum in quadrature	0.041	0.008	0.015

Table 8. Absolute systematic uncertainties in the asymmetry parameter extractions.

of Ξ_c^0 , and no systematic error has been included from this source [48]. Finally, the absolute systematical uncertainties for $\alpha(\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0})$, $\alpha(\Xi_c^0 \rightarrow \Sigma^0\bar{K}^{*0})\alpha(\Sigma^0 \rightarrow \gamma\Lambda)$, and $\alpha(\Xi_c^0 \rightarrow \Sigma^+K^{*-})$ are estimated by adding all individual uncertainties in quadrature. All the uncertainties are summarized in table 8.

7 Summary

We measure for the first time the branching fractions and asymmetry parameters of $\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0}$, $\Xi_c^0 \rightarrow \Sigma^0\bar{K}^{*0}$, and $\Xi_c^0 \rightarrow \Sigma^+K^{*-}$ decays. The relative branching ratios to the normalization mode of $\Xi_c^0 \rightarrow \Xi^-\pi^+$ and the branching fractions of $\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0}$, $\Xi_c^0 \rightarrow \Sigma^0\bar{K}^{*0}$, and $\Xi_c^0 \rightarrow \Sigma^+K^{*-}$ are calculated, as listed in table 3. We note that the branching fraction of $\Xi_c^0 \rightarrow \Sigma^0\bar{K}^{*0}$ is much larger than that of $\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0}$, and this contradicts all the predictions based on $SU(3)_F$ flavor symmetry and dynamical models [7, 24, 26, 31]. This indicates the fraction of $\Lambda K^-\pi^+$ resonating through $\Lambda\bar{K}^{*0}$ is smaller than the fraction of $\Sigma^0 K^-\pi^+$ resonating through $\Sigma^0\bar{K}^{*0}$. The asymmetry parameters $\alpha(\Xi_c^0 \rightarrow \Lambda\bar{K}^{*0})$ and $\alpha(\Xi_c^0 \rightarrow \Sigma^+K^{*-})$ are measured to be $0.15 \pm 0.22(\text{stat.}) \pm 0.04(\text{syst.})$ and $-0.52 \pm 0.30(\text{stat.}) \pm 0.02(\text{syst.})$ with large statistical uncertainties.

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